# Classical Optimal Design on Annulus and Numerical Solution by Shape Derivative Method

### Marko Vrdoljak Joint work with Petar Kunštek



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München, March 2018

State equation ( $\Omega \subseteq \mathbf{R}^d$  open and bounded)

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Two phases:  $0 < \alpha < \beta$  $\mathbf{A} = \chi \alpha \mathbf{I} + (1 - \chi) \beta \mathbf{I}$ ,  $\chi \in L^{\infty}(\Omega; \{0, 1\})$ ,  $\int_{\Omega} \chi \, d\mathbf{x} = q_{\alpha}$ , for given  $0 < q_{\alpha} < |\Omega|$ 

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$$J(\chi) = \int_{\Omega} u(\mathbf{x}) d\mathbf{x} \longrightarrow \max$$

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Interpretations:

- Maximize the amount of heat kept inside body
- Maximize the torsional rigidity of a rod made of two materials
- Maximize the flow rate of two viscous immiscible fluids through pipe

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State equation ( $\Omega \subseteq \mathbf{R}^d$  open and bounded)

$$\left\{ \begin{array}{l} -\mathsf{div}\left(\mathbf{A}\nabla u\right)=1\\ u\in \mathrm{H}_{0}^{1}(\Omega) \end{array} \right.$$

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#### Intuition for annulus?



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### Intuition for annulus?



In general, there might exist no classical optimal design. The relaxation is needed, introducing composite materials.

classical design relaxed design  $\chi \in L^{\infty}(\Omega; \{0, 1\}) \cdots \theta \in L^{\infty}(\Omega; [0, 1])$ 

 $\mathbf{A} \in \mathcal{K}(\theta)$  a.e. on  $\Omega$ 

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State equations

$$\begin{cases} -\operatorname{div} \left( \mathbf{A} \nabla u_i \right) = f_i \\ u_i \in \mathrm{H}^1_0(\Omega) \end{cases} \qquad \qquad i = 1, \dots, m$$

State function  $u = (u_1, \ldots, u_m)$ 

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$$\begin{cases} I(\chi) = \sum_{i=1}^{m} \mu_i \int_{\Omega} f_i u_i \, d\mathbf{x} \to \max\\ \mathsf{u} = (u_1, \dots, u_m) \text{ state function for } \mathbf{A} = \chi \alpha \mathbf{I} + (1 - \chi) \beta \mathbf{I} \\ \chi \in \mathrm{L}^{\infty}(\Omega; \{0, 1\}) \,, \ \int_{\Omega} \chi \, d\mathbf{x} = q_{\alpha} \,, \end{cases}$$

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for some given weights  $\mu_i > 0$ . Relaxed designs:

$$\mathcal{A} := \left\{ (\theta, \mathbf{A}) \in \mathrm{L}^\infty(\Omega; [0, 1] \times \mathrm{M}_d(\mathbf{R})) : \int_\Omega \theta \, d\mathbf{x} = q_\alpha \,, \; \mathbf{A}(\mathbf{x}) \in \mathcal{K}(\theta(\mathbf{x})) \; \text{a.e. on} \; \Omega \right\}$$

Image: A matrix and a matrix

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ight\}$$

However, if  $\Omega$  is a ball and  $f_i$  are radial functions, solution is usually classical. Minimization of the same functional - classical optimal designs are rare exceptions (**Juan Casado-Díaz**); for multiple state problems – joint works with **Krešimir Burazin** and **Ivana Crnjac**.

A. Single state equation [Murat & Tartar, 1985]

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#### A. Single state equation

[**Murat & Tartar, 1985**] There exists relaxed solution  $(\theta^*, \mathbf{A}^*)$  among simple laminates ... conductivity  $\lambda_{\theta}^-$  in one direction  $(\nabla u)$ , and  $\lambda_{\theta}^+$  in orthogonal directions.

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$$rac{1}{\lambda_{ heta}^-} = rac{ heta}{lpha} + rac{1- heta}{eta} \ \lambda_{ heta}^+ = heta lpha + (1- heta)eta$$

can be rewritten as a convex minimization problem

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It is not enough to use only simple laminates, but composite materials that correspond to a non-affine boundary of  $\mathcal{K}(\theta)$  ... higher order sequential laminates. The analogous simpler relaxation fails.

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In spherically symmetric case ( $f_i$  are radial functions), simpler relaxation problem is equivalent to the true relaxation problem (simple laminates are enough).

 $\ldots$  in terms of only local fraction  $\theta$  belonging to the set

$$\mathcal{T}:=\left\{ heta\in\mathrm{L}^\infty(\Omega;[0,1]):\int_\Omega heta\,d\mathbf{x}=q_lpha
ight\}$$

$$\begin{split} I(\theta) &= \sum_{i=1}^{m} \mu_i \int_{\Omega} f_i u_i \, d\mathbf{x} \longrightarrow \max \\ \theta &\in \mathcal{T} \text{ and } u \text{ determined uniquely by} \\ \begin{cases} -\operatorname{div} \left(\lambda_{\theta}^- \nabla u_i\right) = f_i \\ u_i \in \mathrm{H}_0^1(\Omega) \end{cases} \quad i = 1, \dots, m \,, \end{split}$$

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### Minimax formulation

$$\begin{split} I(\theta) &= \sum_{i=1}^{m} \mu_i \int_{\Omega} f_i u_i \, d\mathbf{x} \\ &= -\sum_{i=1}^{m} \mu_i \int_{\Omega} \lambda_{\theta}^{-} |\nabla u_i|^2 - 2f_i u_i \, d\mathbf{x} \\ &= -\min_{\mathbf{v} \in \mathrm{H}_0^1(\Omega; \mathbf{R}^m)} \sum_{i=1}^{m} \mu_i \int_{\Omega} \lambda_{\theta}^{-} |\nabla v_i|^2 - 2f_i v_i \, d\mathbf{x} \\ &= -\max_{\boldsymbol{\sigma} \in \mathcal{S}} \left( -\sum_{i=1}^{m} \mu_i \int_{\Omega} \frac{|\sigma_i|^2}{\lambda_{\theta}^{-}} \, d\mathbf{x} \right) \\ &= \min_{\boldsymbol{\sigma} \in \mathcal{S}} \left( \sum_{i=1}^{m} \mu_i \int_{\Omega} \frac{|\sigma_i|^2}{\lambda_{\theta}^{-}} \, d\mathbf{x} \right) , \end{split}$$

where  $S = \{ \boldsymbol{\sigma} \in L^2(\Omega; \mathbf{R}^d)^m : -\operatorname{div} \sigma_i = f_i, i = 1, \dots, m \}.$ 

### Necessary and sufficient optimality conditions

By minimax theorem there exists a unique  $\sigma^* \in S = \{ \sigma \in L^2(\Omega; \mathbf{R}^d)^m : -\operatorname{div} \sigma_i = f_i, i = 1..m \}$  such that

$$\max_{\mathcal{T}} I = \max_{\theta \in \mathcal{T}} \sum_{i=1}^{m} \mu_i \int_{\Omega} \frac{\beta - \alpha}{\alpha \beta} \, \theta |\boldsymbol{\sigma}_i^*|^2 \, d\mathbf{x} \, .$$

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#### Lemma

The necessary and sufficient condition of optimality for solution  $\theta^* \in \mathcal{T}$  of optimal design problem (1) simplifies to the existence of a Lagrange multiplier  $c \ge 0$  such that

$$\begin{split} &\sum_{\substack{i=1\\m}}^{m} \mu_i |\boldsymbol{\sigma}_i^*|^2 > c \quad \Rightarrow \quad \theta^* = 1 \,, \\ &\sum_{\substack{i=1\\m}}^{m} \mu_i |\boldsymbol{\sigma}_i^*|^2 < c \quad \Rightarrow \quad \theta^* = 0 \,. \end{split}$$

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### Spherically symmetric case – uniqueness

Conditions of optimality:

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In case of spherical symmetry  $\sigma_i^* = \sigma_i^*(r)\mathbf{e}_r$ , where  $\sigma_i^*$  solves  $-\frac{1}{r}(r\sigma_i)' = f_i$ . Let us denote

$$\psi(\mathbf{r}) := \sum_{i=1}^{m} \mu_i |\sigma_i^*|^2 = \sum_{i=1}^{m} \mu_i (\sigma_i^*)^2.$$

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$$\sum_{\substack{i=1\\i=1}}^{m} \mu_i |\boldsymbol{\sigma}_i^*|^2 < c \quad \Rightarrow \quad \theta^* = 0.$$

In case of spherical symmetry  $\sigma_i^* = \sigma_i^*(r)\mathbf{e}_r$ , where  $\sigma_i^*$  solves  $-\frac{1}{r}(r\sigma_i)' = f_i$ . Let us denote

$$\psi(\mathbf{r}) := \sum_{i=1}^{m} \mu_i |\boldsymbol{\sigma}_i^*|^2 = \sum_{i=1}^{m} \mu_i (\sigma_i^*)^2.$$

#### Corollary

For spherically symmetric case, if  $\psi$  is piecewise strictly monotone on  $\omega$  then the problem max<sub>T</sub> I has a unique solution  $\theta^*$ , which is a radial characteristic function. Consequently, the solution of the true relaxation problem is unique, classical and radial.

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### Optimal design problem on annulus



Single state equation:

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where 
$$\lambda_{\theta(\mathbf{x})}^{-} = \left(\frac{\theta(\mathbf{x})}{\alpha} + \frac{1-\theta(\mathbf{x})}{\beta}\right)^{-1}$$
.

**Optimization problem:** 

$$\begin{cases} I(\theta) = \int_{\Omega} u \, \mathrm{d} \mathbf{x} \to \max \\ s.t. \quad \theta \in \mathrm{L}^{\infty}(\Omega, [0, 1]), \quad \int_{\Omega} \theta = q_{\alpha}, \text{ where } u \text{ satisfies (2)} \end{cases}$$
(3)

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(2)

One can rewrite (2) in polar coordinates :

$$-\frac{1}{r^{d-1}}(r^{d-1}\underbrace{\lambda_{\theta}^{-}u'(r)}_{\sigma})'=1 \text{ in } \langle r_1,r_2\rangle, \quad u(r_1)=u(r_2)=0.$$

Observe that  $\sigma$  satisfies

$$\sigma = -\frac{r}{d} + \frac{\gamma}{r^{d-1}}, \quad \gamma > 0$$

 $\sigma(\mathbf{r}): \langle \mathbf{0}, \infty \rangle \to \mathbb{R}$  is a strictly decreasing function, for any  $\gamma$ .

 $\implies$  Optimal design is unique, classical and radial.

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The necessary and sufficient condition of optimality for  $\theta^*$  states



$$egin{array}{ccc} |\sigma^*| > c &\Rightarrow& heta^* = 1\,, \ |\sigma^*| < c &\Rightarrow& heta^* = 0\,. \end{array}$$

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$$egin{array}{ccc} |\sigma^*| > c &\Rightarrow& heta^* = 1\,, \ |\sigma^*| < c &\Rightarrow& heta^* = 0\,. \end{array}$$

There are only three possible candidates for optimal design:

1) 
$$\theta^{*}(r) = \begin{cases} 1, & r \in [r_{1}, r_{+}) \\ 0, & r \in [r_{+}, r_{-}) \\ 1, & r \in [r_{-}, r_{2}] \end{cases}$$
  
2)  $\theta^{*}(r) = \begin{cases} 1, & r \in [r_{1}, r_{+}) \\ 0, & r \in [r_{+}, r_{2}) \end{cases}$   
3)  $\theta^{*}(r) = \begin{cases} 0, & r \in [r_{1}, r_{-}) \\ 1, & r \in [r_{-}, r_{2}) \end{cases}$ 

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From condition of optimality a non-linear system (with unknowns  $\gamma$ , c,  $r_+$ ,  $r_-$ ) is created:

$$\begin{cases} S_d \int_{r_1}^{r_2} \theta(\rho) \rho^{d-1} d\rho = q_\alpha \\ u(r_2) = 0 \iff \gamma \int_{r_1}^{r_2} \left(\frac{1}{a(\rho)\rho^{d-1}}\right) d\rho = \int_{r_1}^{r_2} \frac{\rho}{a(\rho)} d\rho \\ \sigma(r_+) = c, \quad \sigma(r_-) = -c, \quad \text{where } c > 0 \end{cases}$$
(NS)

where

$$\sigma(r) = \frac{\gamma}{r^{d-1}} - \frac{r}{d}, \quad \& \quad a(r) = \left(\frac{\theta(r)}{\alpha} + \frac{1 - \theta(r)}{\beta}\right)^{-1}$$

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#### Theorem (Optimal design for annulus d = 2, 3, f = 1)

With previous assumptions the problem admits classical solution with two possible designs:

eta-alpha

1) 
$$\theta^{*}(r) = \begin{cases} 1, & r \in [r_{1}, r_{+}) \\ 0, & r \in [r_{+}, r_{-}) \\ 1, & r \in [r_{-}, r_{2}] \end{cases}$$
 alpha-beta-alpha-beta-alpha-beta-alpha-beta alpha-beta alpha-beta

More precisely, if  $q_{\alpha}$  is small enough, design 2) is optimal.



### Shape derivative

Perturbation of the set  $\boldsymbol{\Omega}$  is given with

$$\Omega_t = (\mathsf{Id} + t\psi)\Omega$$

where  $\psi \in W^{k,\infty}(\mathbf{R}^d, \mathbf{R}^d)$ .



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Perturbation of the set  $\boldsymbol{\Omega}$  is given with

$$\Omega_t = (\mathsf{Id} + t\psi)\Omega$$

where  $\psi \in W^{k,\infty}(\mathbf{R}^d, \mathbf{R}^d)$ .



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#### Definition (Shape derivative)

Let  $J = J(\Omega)$  be a shape functional. J is said to be shape differentiable at  $\Omega$  in direction  $\psi$  if

$$J'(\Omega,\psi):=\lim_{t\searrow 0}rac{J(\Omega_t)-J(\Omega)}{t}$$

exists and the mapping  $\psi \mapsto J'(\Omega, \psi)$  is linear and continuous.  $J'(\Omega, \psi)$  is called the **shape derivative**.

### Single state problem (general f)

In case of our single state optimal design problem:

$$J(\Omega_{lpha}) = \int_{\Omega} \mathit{fu} \, \mathrm{d} \mathbf{x} o \mathsf{max}$$

u is determined by  $\mathbf{A} = \chi \alpha \mathbf{I} + (1 - \chi) \beta \mathbf{I}$  $\chi \in L^{\infty}(\Omega, \{0, 1\})$  is a characteristic function of  $\Omega_{\alpha}$ ,  $|\Omega_{\alpha}| = q_{\alpha}$ 

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### Single state problem (general f)

In case of our single state optimal design problem:

$$J(\Omega_{lpha}) = \int_{\Omega} f u \, \mathrm{d} \mathbf{x} o \mathsf{max}$$

u is determined by  $\mathbf{A} = \chi \alpha \mathbf{I} + (1 - \chi) \beta \mathbf{I}$  $\chi \in L^{\infty}(\Omega, \{0, 1\})$  is a characteristic function of  $\Omega_{\alpha}$ ,  $|\Omega_{\alpha}| = q_{\alpha}$ 

The shape derivative is given by:

$$\begin{split} J'(\Omega_{\alpha},\psi) &= \int_{\Omega} \mathbf{A}(-\mathsf{div}\,(\psi) + \nabla\psi + \nabla\psi^{\tau}) \nabla u_{0} \cdot \nabla u_{0} \,\mathrm{d}\mathbf{x} \\ &+ \int_{\Omega} 2(\mathsf{div}\,(\psi)f + \nabla f \cdot \psi) u_{0} \,\mathrm{d}\mathbf{x} \end{split}$$

where  $u_0$  is the corresponding state function.

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### Single state problem (general f)

In case of our single state optimal design problem:

$$J(\Omega_{lpha}) = \int_{\Omega} f u \, \mathrm{d} \mathbf{x} o \max$$

u is determined by  $\mathbf{A} = \chi \alpha \mathbf{I} + (1 - \chi) \beta \mathbf{I}$  $\chi \in L^{\infty}(\Omega, \{0, 1\})$  is a characteristic function of  $\Omega_{\alpha}$ ,  $|\Omega_{\alpha}| = q_{\alpha}$ 

The shape derivative is given by:

$$\begin{split} J'(\Omega_{\alpha},\psi) &= \int_{\Omega} \mathbf{A}(-\mathsf{div}\,(\psi) + \nabla\psi + \nabla\psi^{\tau}) \nabla u_0 \cdot \nabla u_0 \,\mathrm{d}\mathbf{x} \\ &+ \int_{\Omega} 2(\mathsf{div}\,(\psi)f + \nabla f \cdot \psi) u_0 \,\mathrm{d}\mathbf{x} \end{split}$$

where  $u_0$  is the corresponding state function. The construction of  $\psi$ :

$$\int_{\Omega} \nabla \psi : \nabla \varphi + \int_{\Omega} \psi \cdot \varphi = \mathcal{L}'(\Omega_{\alpha}, \varphi), \quad \forall \varphi \in H^1_0(\Omega).$$

(4)

#### Algorithm 1: k-th step of gradient method

- 1 Input :  $\Gamma_k$  interface is discretized in points (it is used to create new mesh  $\mathcal{T}_k$ )
- 2 Construct vector spaces Vh on mesh  $\mathcal{T}_k$  (Vh=P1, P2 ...)
- 3 Determine vector field  $\psi \in \mathtt{Vh}$
- 4 Determine  $t_0 > 0$  (if too small, increase of J is insignificant; upper bound is dictated by mesh  $T_k$ )
- 5 Move mesh:  $\mathcal{T}_{k+1} = (\mathsf{Id} + t_0 \psi) \mathcal{T}_k$
- 6 Output:  $\Gamma_{k+1} = (\mathsf{Id} + t_0 \psi) \Gamma_k$

Upper bound in part 4 is calculated by checkmovemesh. This ensures that moving of a mesh doesn't create wrong ordering of elements (volume of triangle should not be negative). Part 5 is implemented using movemesh. At the end of the step it is recommended to use adaptmesh.

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# Thank you for your attention!

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