## Explicit solutions of multiple state optimal design problems

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Joint work with Marko Vrdoljak


## Outline

(1) Energy minimization and relaxation

Posing the problem
Relaxation
(2) Convex minimization problem

Simpler problem
Spherically symmetric case
(3) Examples

One state
Multiple states

## Multiple state optimal design problem

$\Omega \subseteq \mathbf{R}^{d}$ open and bounded, $f_{1}, \ldots, f_{m} \in \mathrm{~L}^{2}(\Omega)$ given; stationary diffusion equations with homogenous Dirichlet b. c.:

$$
\left\{\begin{array}{l}
-\operatorname{div}\left(\mathbf{A} \nabla u_{i}\right)=f_{i}  \tag{1}\\
u_{i} \in \mathrm{H}_{0}^{1}(\Omega)
\end{array}, \quad i=1, \ldots, m\right.
$$

where $\mathbf{A}$ is a mixture of two isotropic materials with conductivities $0<\alpha<\beta: \mathbf{A}=\chi \alpha \mathbf{I}+(1-\chi) \beta \mathbf{I}$, where $\chi \in \mathrm{L}^{\infty}(\Omega ;\{0,1\})$, For given $\Omega, \alpha, \beta, q_{\alpha}, f_{i}$, and some given weights $\mu_{i}>0$, we want to find such material $\mathbf{A}$ which minimizes the weighted sum of energies (total amounts of heat/electrical energy dissipated in $\Omega$ ):


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$$
I(\chi):=\sum_{i=1}^{m} \mu_{i} \int_{\Omega} f_{i} u_{i} d \mathbf{x} \rightarrow \min , \quad \chi \in \mathrm{~L}^{\infty}(\Omega ;\{0,1\})
$$

Murat \& Tartar

## Lurie \& Cherkaev



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\text { classical material }
\end{array}
$$

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$$
\begin{array}{rll}
\chi \in \mathrm{L}^{\infty}(\Omega ;\{0,1\}) & \cdots & \theta \in \mathrm{L}^{\infty}(\Omega ;[0,1]) \\
\mathbf{A}=\chi \alpha \mathbf{I}+(1-\chi) \beta \mathbf{I} & & \mathbf{A} \in \mathcal{K}(\theta) \text { a.e. on } \Omega \\
\text { classical material } & & \text { composite mateiral - relaxation }
\end{array}
$$

## Composite material

## Definition

If a sequence of characteristic functions $\chi_{\varepsilon} \in \mathrm{L}^{\infty}(\Omega ;\{0,1\})$ and conductivities $\mathbf{A}^{\varepsilon}(x)=\chi_{\varepsilon}(x) \alpha \mathbf{I}+\left(1-\chi_{\varepsilon}(x)\right) \beta \mathbf{I}$ satisfy $\chi_{\varepsilon} \rightharpoonup \theta$ weakly * and $\mathbf{A}^{\varepsilon} H$-converges to $\mathbf{A}^{*}$, then it is said that $\mathbf{A}^{*}$ is homogenised tensor of two-phase composite material with proportions $\theta$ of first material and microstructure defined by the sequence $\left(\chi_{\varepsilon}\right)$.

Example - simple laminates: if $\chi_{\varepsilon}$ depend only on $x_{1}$, then
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Set of all composites:

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\mathbf{A}^{*}=\operatorname{diag}\left(\lambda_{\theta}^{-}, \lambda_{\theta}^{+}, \lambda_{\theta}^{+}, \ldots, \lambda_{\theta}^{+}\right),
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where

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\lambda_{\theta}^{+}=\theta \alpha+(1-\theta) \beta, \quad \frac{1}{\lambda_{\theta}^{-}}=\frac{\theta}{\alpha}+\frac{1-\theta}{\beta}
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Set of all composites:

$$
\mathcal{A}:=\left\{(\theta, \mathbf{A}) \in \mathrm{L}^{\infty}\left(\Omega ;[0,1] \times \mathrm{M}_{d}(\mathbf{R})\right): \int_{\Omega} \theta d \mathbf{x}=q_{\alpha}, \mathbf{A} \in \mathcal{K}(\theta) \text { a.e. }\right\}
$$

Effective conductivities - set $\mathcal{K}(\theta)$
G-closure problem: for given $\theta$ find all possible homogenised (effective) tensors $\mathbf{A}^{*}$
$\mathcal{K}(\theta)$ is given in terms of eigenvalues
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2D:


$$
\begin{aligned}
& \sum_{j=1}^{d} \frac{1}{\lambda_{j}-\alpha} \leq \frac{1}{\lambda_{\theta}^{-}-\alpha}+\frac{d-1}{\lambda_{\theta}^{+}-\alpha} \\
& \sum_{j=1}^{d} \frac{1}{\beta-\lambda_{j}} \leq \frac{1}{\beta-\lambda_{\theta}^{-}}+\frac{d-1}{\beta-\lambda_{\theta}^{+}}
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\min _{L^{\infty}(\Omega ;\{0,1\})} I
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## How do we find a solution?

Goal: find explicit solution for some simple domains (circle)
Motivation: test examples for robust numerical algorithms
A. Single state equation: [Murat \& Tartar] This problem can be rewritten as a simpler convex minimization problem.

$\mathcal{T}=\left\{\theta \in \mathrm{L}^{\infty}(\Omega ;[0,1]): \int_{\Omega} \theta=q_{\alpha}\right\}$
$\theta \in \mathcal{T}$. and $u$ determined uniquely by

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\left\{\begin{array}{l}
-\operatorname{div}\left(\lambda_{\theta}^{+} \nabla u\right)=f \\
u \in \mathrm{H}_{0}^{1}(\Omega)
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B. Multiple state equations: Simpler
relaxation fails; in spherically symmetric
case or when $m<d$, it can be done!

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I(\theta)=\sum_{i=1}^{m} \mu_{i} \int_{\Omega} f_{i} u_{i} d x \rightarrow \min
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\end{array}
$$

## Theorem

If $m<d$ then $\min _{\mathcal{A}} J=\min _{\mathcal{T}} I$ and:

- There is unique $\mathrm{u}^{*} \in \mathrm{H}_{0}^{1}\left(\Omega ; \mathbf{R}^{m}\right)$ which is the state for every solution of $\min _{\mathcal{A}} J$ and $\min _{\mathcal{T}} I$.
- If $\left(\theta^{*}, \mathbf{A}^{*}\right)$ is an optimal design for the problem $\min _{\mathcal{A}} J$, then $\theta^{*}$ is optimal design for $\min _{\mathcal{T}} I$.
- Conversely, if $\theta^{*}$ is a solution of optimal design problem min $\mathcal{T}$ I, then any $\left(\theta^{*}, \mathbf{A}^{*}\right) \in \mathcal{A}$ satisfying $\mathbf{A}^{*} \nabla u_{i}^{*}=\lambda_{\theta^{*}}^{+} \nabla u_{i}^{*}$, $i=1, \ldots, m$ almost everywhere on $\Omega$ (e.g. simple laminates) is an optimal design for the problem $\min _{\mathcal{A}} J$.


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$i=1, \ldots, m$ almost everywhere on $\Omega$ (e.g. simple laminates) is an optimal design for the problem $\min _{\mathcal{A}} J$.

Spherical symmetry: $\min _{\mathcal{A}} J \Longleftrightarrow \min _{\mathcal{T}} I$

## Theorem

Let $\Omega \subseteq \mathbf{R}^{d}$ be spherically symmetric, and let the right-hand sides $f_{i}=f_{i}(r), r \in \omega, i=1, \ldots, m$ be radial functions. Then $\min _{\mathcal{A}} J=\min _{\mathcal{T}} I$ and there is unique (radial) $\mathrm{u}^{*}$ which is the state for any solution of $\min _{\mathcal{A}} J$ and $\min _{\mathcal{T}} I$. Moreover,

where $S$ denotes the surface measure on a sphere. Then $\theta^{*}$ is also minimizer for $I$ over $\mathcal{T}$.

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a) For any minimizer $\theta$ of functional $I$ over $\mathcal{T}$, let us define a radial function $\theta^{*}: \Omega \longrightarrow \mathbf{R}$ as the average value over spheres of $\theta$ : for $r \in \omega$ we take

$$
\theta^{*}(r):=f_{\partial B(\mathbf{0}, r)} \theta d S
$$

where $S$ denotes the surface measure on a sphere. Then $\theta^{*}$ is also minimizer for I over $\mathcal{T}$.

## Spherical symmetry...cont.

## Theorem

b) For any radial minimizer $\theta^{*}$ of $I$ over $\mathcal{T}$, let us define $\mathbf{A}^{*} \in \mathcal{K}\left(\theta^{*}\right)$ as a simple laminate with the lamination direction orthogonal to the radial vector $\mathbf{e}_{r}$, almost everywhere on $\Omega$. To be specific, we define


## Spherical symmetry... cont.

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\mathbf{A}^{*}(\mathbf{x}):=\operatorname{diag}\left(\lambda_{\theta^{*}}^{+}(|\mathbf{x}|), \lambda_{\theta^{*}}^{-}(|\mathbf{x}|), \lambda_{\theta^{*}}^{+}(|\mathbf{x}|), \ldots, \lambda_{\theta^{*}}^{+}(|\mathbf{x}|)\right)
$$

in spherical basis $\left(\mathbf{e}_{r}(\mathbf{x}), \mathbf{e}_{\phi_{1}}(\mathbf{x}), \mathbf{e}_{\phi_{2}}(\mathbf{x}), \ldots, \mathbf{e}_{\phi_{d-1}}(\mathbf{x})\right)$. Then $\left(\theta^{*}, \mathbf{A}^{*}\right)$ is an optimal design for $\min _{\mathcal{A}} J$.
 $i=1$

## Spherical symmetry... cont.

## Theorem

b) For any radial minimizer $\theta^{*}$ of $I$ over $\mathcal{T}$, let us define $\mathbf{A}^{*} \in \mathcal{K}\left(\theta^{*}\right)$ as a simple laminate with the lamination direction orthogonal to the radial vector $\mathbf{e}_{r}$, almost everywhere on $\Omega$. To be specific, we define

$$
\mathbf{A}^{*}(\mathbf{x}):=\operatorname{diag}\left(\lambda_{\theta^{*}}^{+}(|\mathbf{x}|), \lambda_{\theta^{*}}^{-}(|\mathbf{x}|), \lambda_{\theta^{*}}^{+}(|\mathbf{x}|), \ldots, \lambda_{\theta^{*}}^{+}(|\mathbf{x}|)\right)
$$

in spherical basis $\left(\mathbf{e}_{r}(\mathbf{x}), \mathbf{e}_{\phi_{1}}(\mathbf{x}), \mathbf{e}_{\phi_{2}}(\mathbf{x}), \ldots, \mathbf{e}_{\phi_{d-1}}(\mathbf{x})\right)$. Then $\left(\theta^{*}, \mathbf{A}^{*}\right)$ is an optimal design for $\min _{\mathcal{A}} J$.
c) If $\left(\theta^{*}, \mathbf{A}^{*}\right) \in \mathcal{A}$ is a solution of the relaxed problem $\min _{\mathcal{A}} J$ then $\theta^{*}$ is optimal for $\min _{\mathcal{T}} I$, and $\mathbf{A}^{*} \nabla u_{i}^{*}=\lambda_{\theta^{*}}^{+} \nabla u_{i}^{*}$, almost everywhere, $i=1, \ldots, m$.

## Uniqueness on a ball

## Lemma

Let $\Omega$ be ball $B(\mathbf{0}, R)$, and let the right-hand sides $f_{i}$ be radial functions, such that mappings $r \mapsto r^{\frac{d-1}{2}} f_{i}(r)$ belong to $\mathrm{L}^{2}(\langle 0, R\rangle), i=1, \ldots, m$. Then there are unique radial fluxes

$$
\sigma_{i}^{*}(r)=-\frac{1}{r^{d-1}} \int_{0}^{r} \rho^{d-1} f_{i}(\rho) d \rho \mathbf{e}_{r}
$$

corresponding to each minimizer of $\min _{\mathcal{T}} I$, and this minimizer is radial and unique on the set where at least one $\sigma_{i}^{*}$ does not vanish. If the Lagrange multiplier c is positive, this holds true on the whole $B(\mathbf{0}, R)$.

## Optimality conditions for $\min _{\mathcal{T}} I$

## Lemma

$\theta^{*} \in \mathcal{T}$ is a solution $\min \mathcal{T} I$ if and only if there exists a Lagrange multiplier $c \geq 0$ such that
or equivalently

$$
\begin{aligned}
\theta^{*} \in\langle 0,1\rangle & \Rightarrow \sum_{i=1}^{m} \mu_{i}\left|\nabla u_{i}^{*}\right|^{2}=c, \\
\theta^{*}=0 & \Rightarrow \sum_{i=1}^{m} \mu_{i}\left|\nabla u_{i}^{*}\right|^{2} \geq c, \\
\theta^{*}=1 & \Rightarrow \sum_{i=1}^{m} \mu_{i}\left|\nabla u_{i}^{*}\right|^{2} \leq c,
\end{aligned}
$$

$$
\begin{aligned}
& \sum_{i=1}^{m} \mu_{i}\left|\nabla u_{i}^{*}\right|^{2}>c \Rightarrow \theta^{*}=0, \\
& \sum_{i=1}^{m} \mu_{i}\left|\nabla u_{i}^{*}\right|^{2}<c \Rightarrow \theta^{*}=1 .
\end{aligned}
$$

## Ball $\Omega=B(\mathbf{0}, 2) \subseteq \mathbf{R}^{2}$ with nonconstant right-hand side

In all examples $\alpha=1, \beta=2$, one state equation, $f(r)=1-r$

## State equation in polar coordinates



Conditions of optimality: there exists a constant $\gamma:=\sqrt{c}>0$ such that for optimal $\theta^{*}$ we have:


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-\frac{1}{r}\left(r \lambda_{\theta(r)}^{+} u^{\prime}\right)^{\prime}=1-r .
$$ Integration gives $\quad\left|u^{\prime}(r)\right|=\frac{\psi(r)}{\alpha \theta(r)+\beta(1-\theta(r))}, \quad$ where $\psi(r)=\frac{\left|2 r^{2}-3 r\right|}{6}$.

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& \Rightarrow g_{\beta}:=\frac{\psi}{\beta}>\gamma \\
\left|u^{\prime}(r)\right|<\gamma & \Rightarrow \theta^{*}(r)=1 \\
& \Rightarrow g_{\alpha}:=\frac{\psi}{\alpha}<\gamma \\
\theta^{*} \in\langle 0,1\rangle & \Rightarrow\left|u^{\prime}(r)\right|=\gamma \\
& \Rightarrow \theta^{*}(r)=\frac{\beta \gamma-\psi(r)}{\gamma(\beta-\alpha)}
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Lagrange multiplier $\gamma$ is uniquely determined by the constraint $f_{\Omega} \theta^{*} d \mathbf{x}=\eta:=\frac{q_{\alpha}}{|\Omega|} \in[0,1]$, which is algebraic equation for $\gamma$.

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Two state equations on a ball $\Omega=B(0,2)$

- $f_{1}=\chi_{B(\mathbf{0}, 1)}, f_{2} \equiv 1$,
- $\left\{-\operatorname{div}\left(\lambda_{\theta}^{+} \nabla u_{i}\right)=f_{i}\right.$ $i=1,2$
- $\mu \int_{\Omega} f_{1} u_{1} d \mathbf{x}+\int_{\Omega} f_{2} u_{2} d \mathbf{x} \rightarrow \min$


## Solving state equation


with


Similarly as in the first example: $\psi:=\mu \psi_{1}^{2}+\psi_{2}^{2}, g_{\alpha}:=\frac{\psi}{\alpha^{2}}, g_{\beta}:=\frac{\psi}{\beta^{2}}$.

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Solving state equation

$$
u_{i}^{\prime}(r)=\frac{\psi_{i}(r)}{\theta(r) \alpha+(1-\theta(r)) \beta}, i=1,2,
$$

with

$$
\psi_{1}(r)=\left\{\begin{aligned}
-\frac{r}{2}, & 0 \leq r<1, \\
-\frac{1}{2 r}, & 1 \leq r \leq 2
\end{aligned} \quad \text { and } \psi_{2}(r)=-\frac{r}{2}\right.
$$

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## Geometric interpretation of optimality conditions



As before, Lagrange multiplier can be numerically calculated from corresponding algebraic equation $f_{\Omega} \theta^{*} d \mathrm{x}=\eta$.

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