

Complex Friedrichs systems and applications

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Abstract settings

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- L complex Hilbert space (L' antidual of L),
- $\mathcal{D} \subseteq L$ dense subspace
- $\mathcal{L}, \tilde{\mathcal{L}}: \mathcal{D} \to L$ linear unbounded operators satisfying

$$(\forall \varphi, \psi \in \mathcal{D}) \qquad \langle \mathcal{L}\varphi | \psi \rangle_L = \langle \varphi | \tilde{\mathcal{L}}\psi \rangle_L, \qquad (T1)$$

$$(\exists c > 0) (\forall \varphi \in \mathcal{D}) \qquad \| (\mathcal{L} + \tilde{\mathcal{L}}) \varphi \|_L \le c \| \varphi \|_L, \tag{T2}$$

$$(\exists \mu_0 > 0) (\forall \varphi \in \mathcal{D}) \quad \langle (\mathcal{L} + \tilde{\mathcal{L}}) \varphi | \varphi \rangle_L \ge 2\mu_0 \|\varphi\|_L^2. \tag{T3}$$

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Classical complex Freidrichs operator

Let $d, r \in \mathbf{N}, \Omega \subseteq \mathbf{R}^d$ open and bounded with Lipshitz boundary, $\mathcal{D} = C_c^{\infty}(\Omega; \mathbf{C}^r), L = L^2(\Omega; \mathbf{C}^r), \mathbf{A}_k \in \mathrm{W}^{1,\infty}(\Omega; \mathrm{M}_r(\mathbf{C})), k \in 1..d$ and $\mathbf{C} \in \mathrm{L}^{\infty}(\Omega; \mathrm{M}_r(\mathbf{C}))$ satisfying

Operators $\mathcal{L}, \tilde{\mathcal{L}}: \mathcal{D} \to L$ defined as

$$\begin{split} \mathcal{L}\mathbf{u} &:= \quad \sum_{k=1}^{d} \partial_k(\mathbf{A}_k\mathbf{u}) + \mathbf{C}\mathbf{u} \\ \tilde{\mathcal{L}}\mathbf{u} &:= \quad -\sum_{k=1}^{d} \partial_k(\mathbf{A}_k^*\mathbf{u}) + (\mathbf{C}^* + \sum_{k=1}^{d} \partial_k\mathbf{A}_k^*)\mathbf{u} \end{split}$$

satisfy (T1)-(T3).

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Operator $\ensuremath{\mathcal{L}}$ is called the symmetric positive operator or the Friedrichs operator and

$$\mathcal{L}\mathsf{u}=\mathsf{f}$$

the symmetric positive system or the Friedrichs system.

- introduced in K. O. Friedrichs: Symmetric positive linear differential equations, Communications on Pure and Applied Mathematics 11 (1958), 333-418
- goal: treating the equations of mixed type, such as the Tricommi equation

$$y\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$

unified tretment of equations and systems of different type

- convenient for the numerical treatment

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Formulation of the problem

• $(\mathcal{D}, \langle \,\cdot\, |\,\cdot\, \rangle_{\mathcal{L}})$ is an inner product space, where

$$\langle \, \cdot \, | \, \cdot \, \rangle_{\mathcal{L}} := \langle \, \cdot \, | \, \cdot \, \rangle_L + \langle \, \mathcal{L} \, \cdot \, | \, \mathcal{L} \cdot \, \rangle_L \, .$$

- $\|\cdot\|_{\mathcal{L}}$ is called the graph norm.
- W_0 the completion of $\mathcal D$ in the graph norm $\dots \mathcal L, \tilde{\mathcal L} \in \mathcal L(L; W_0')$

Lemma

The graph space

$$W := \{ u \in L : \mathcal{L}u \in L \} = \{ u \in L : \tilde{\mathcal{L}}u \in L \}$$

is Hilbert space with respect to $\langle \cdot | \cdot \rangle_{\mathcal{L}}$.

Problem: for given $f \in L$ find $u \in W$ such that $\mathcal{L}u = f$.

Find sufficient conditions on $V \leq W$ such that $\mathcal{L}_{|V} : V \to L$ is an isomorphism.

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Boundary operator

Boundary operator $D \in \mathcal{L}(W;W')$ is defined by

$$_{W'}\langle Du,v\rangle_W := \langle \mathcal{L}u \,|\, v \,\rangle_L - \langle \, u \,|\, \tilde{\mathcal{L}}v \,\rangle_L \qquad u, \, v \in W.$$

Lemma

Under assumptions (T1)–(T2), operator D is selfadjoint

$$_{W'}\langle Du,v\rangle_W = \overline{_{W'}\langle Dv,u\rangle_W}$$

and satisfies

$$\ker D = W_0$$

$$\operatorname{im} D = W_0^0 := \{g \in W' : (\forall u \in W_0) \mid_{W'} \langle g, u \rangle_W = 0\}.$$

In particular, $\operatorname{im} D$ is closed in W'.

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Well-posedness theorem

Let V and \tilde{V} be subspaces of W that satisfy

$$\begin{array}{ll} (\forall u \in V) & _{W'} \langle Du, u \rangle_W \geq 0 \\ (\forall v \in \tilde{V}) & _{W'} \langle Dv, v \rangle_W \leq 0 \end{array}$$
 (V1)

$$V = D(\tilde{V})^o, \qquad \tilde{V} = D(V)^o. \tag{V2}$$

Theorem

Under assumptions (T1)–(T3) and (V1)–(V2), the operators $\mathcal{L}_{|V}: V \to L$ and $\tilde{\mathcal{L}}_{|\tilde{V}}\tilde{V} \to L$ are isomorphisms.¹

¹In real case: [AE&JLG&GC2007].



Non-stationary complex Friedrichs systems

Consider abstract Cauchy problem

$$\begin{cases} \mathsf{u}'(t) + \mathcal{L}\mathsf{u}(t) = \mathsf{f} \\ \mathsf{u}(0) = \mathsf{u}_0, \end{cases}$$

where $u: [0, T\rangle \rightarrow L, T > 0$ is the unknown function, $f: \langle 0, T\rangle \rightarrow L$, $u_0 \in L$ and \mathcal{L} is abstract Friedrichs operator that satisfy (T1)–(T2) and

$$(\forall \varphi \in \mathcal{D}) \qquad \operatorname{Re} \left\langle (\mathcal{L} + \tilde{\mathcal{L}}) \varphi \, | \, \varphi \right\rangle_L \ge 0.$$
 (T3')

Let $V \leq W$ satisfy (V1)–(V2). Then following is valid

Theorem

 $-\mathcal{L}_{|V}$ is the infinitesimal generator of a contraction C_0 -semigroup $(T(t))_{t\geq 0}$ on L^2

²In real case: [BE2016].



Theorem

Let \mathcal{L} be an operator that satisfy (T1)–(T2) and (T3'), V subspace of its graph space satisfying (V1)–(V2), and $f \in L^1(\langle 0, T \rangle; L)$. Then for every $u_0 \in L$ Cauchy problem

$$\begin{cases} \mathsf{u}'(t) + \mathcal{L}\mathsf{u}(t) = \mathsf{f} \\ \mathsf{u}(0) = \mathsf{u}_0 \end{cases}$$

has the unique weak solution $\mathbf{u} \in C([0,T];L)$ given by

$$u(t) = T(t)u_0 + \int_0^t T(t-s)f(s) \, ds, \quad t \in [0,T],$$

where $(T(t))_{t\geq 0}$ is a contraction C_0 -semigroup generated with $-\mathcal{L}_{|V}$.

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Friedrichs systems in H^s spaces

Let $s \in \mathbf{R}$, $L = \mathrm{H}^{s}(\mathbf{R}^{d}; \mathbf{C}^{r})$, $\mathcal{D} = \mathrm{C}^{\infty}_{c}(\mathbf{R}^{d}; \mathbf{C}^{r})$ and assume that constant matrices \mathbf{C} , \mathbf{A}_{k} , $k \in 1..d$, satisfy (F1) and (F2):

$$\mathbf{A}_k = \mathbf{A}_k^*$$
,

$$(\exists \mu_0 > 0) \quad \mathbf{C} + \mathbf{C}^* \ge 2\mu_0 \mathbf{I} \qquad (\text{ae on } \Omega) \;.$$

Operators

$$\mathcal{L} u := \sum_{k=1}^d \partial_k (\mathbf{A}_k u) + \mathbf{C} u$$

and

$$\tilde{\mathcal{L}} \textbf{u} := -\sum_{k=1}^d \partial_k (\mathbf{A}_k \textbf{u}) + \mathbf{C}^* \textbf{u}$$

satisfy (T1)–(T3), boundary operator D is trival and V = V = W.

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Linear Dirac system

Consider system of equations

$$\gamma^{0}\partial_{t}\psi + \gamma^{1}\partial_{1}\psi + \gamma^{2}\partial_{2}\psi + \gamma^{3}\partial_{3}\psi + \mathbf{B}\psi = \mathbf{f}, \tag{1}$$

where $\psi : [0, T \rangle \times \mathbf{R}^3 \to \mathbf{C}^4$ is unknown function, $f : \langle 0, T \rangle \to \mathbf{C}^4$, $\mathbf{B} = \begin{bmatrix} b_1 \mathbf{I} & 0 \\ 0 & b_2 \mathbf{I} \end{bmatrix}$ for $b_1, b_2 : \mathbf{R}^3 \to \mathbf{C}$ and 2×2 unit matrix \mathbf{I} , while

$$\gamma^0 = \begin{bmatrix} \mathbf{I} & 0\\ 0 & -\mathbf{I} \end{bmatrix}, \qquad \gamma^k = \begin{bmatrix} 0 & \sigma^k\\ -\sigma^k & 0 \end{bmatrix}, \quad k = 1, 2, 3.$$

$$\sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \qquad \sigma^2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \qquad \sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

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System (1) can be written as Friedrichs system

$$\partial_t \psi + \mathcal{L} \psi = \mathsf{F},$$

where
$$\mathsf{F} = \gamma^0 \mathsf{f}$$
 and $\mathcal{L}\psi = \sum_{k=1}^3 \mathbf{A}_k \partial_k \psi + \mathbf{C}\psi$ for $\mathbf{A}_k = \tilde{\mathbf{A}}_k := \begin{bmatrix} 0 & \sigma^k \\ \sigma^k & 0 \end{bmatrix}, \qquad \mathbf{C} = \gamma^0 \mathbf{B}.$

Spaces involved:

$$\begin{aligned} \mathcal{D} &= \mathrm{C}_c^{\infty}(\mathbf{R}^3; \mathbf{C}^4) \\ L &= \mathrm{L}^2(\mathbf{R}^3; \mathbf{C}^4), (\text{or } \mathrm{H}^s(\mathbf{R}^3; \mathbf{C}^4)) \\ W &= \{ \mathsf{u} \in \mathrm{L}^2(\mathbf{R}^3; \mathbf{C}^4) : \sum_{k=1}^3 \mathbf{A}_k \partial_k \mathsf{u} \in \mathrm{L}^2(\mathbf{R}^3; \mathbf{C}^4) \} \end{aligned}$$

D is trivial, $V=\tilde{V}=W.$

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Dirac-Klein-Gordon system

$$\begin{cases} -i(\gamma^0\partial_t + \gamma^1\partial_1 + \gamma^2\partial_2 + \gamma^3\partial_3 + M)\psi = \phi\psi\\ \partial_t^2\phi - \Delta\phi + m^2\phi = \psi^*\gamma^0\psi \end{cases}$$
(2)

where unknown functions are $\psi = \psi(t, x) : \mathbf{R}^{1+3} \to \mathbf{C}^4$ and $\phi : \mathbf{R}^{1+3} \to \mathbf{R}$, while $M, m \ge 0$ and $\gamma^k, k = 1..3$ are same as in previous example.

Remark

For two Friedrichs systems

$$\partial_t \mathbf{u}_1 + \mathcal{L}_1 \mathbf{u}_1 = \mathbf{f}_1 \partial_t \mathbf{u}_2 + \mathcal{L}_2 \mathbf{u}_2 = \mathbf{f}_2$$

system

$$\begin{array}{l} \partial_t \mathsf{u} + \mathcal{L} \mathsf{u} = \mathsf{f} \\ \text{s also a Friedrichs system with } \mathcal{L} = \begin{bmatrix} \mathcal{L}_1 & 0 \\ 0 & \mathcal{L}_2 \end{bmatrix}, \, \mathsf{u} = \begin{bmatrix} \mathsf{u}_1 \\ \mathsf{u}_2 \end{bmatrix}, \mathsf{f} = \begin{bmatrix} \mathsf{f}_1 \\ \mathsf{f}_2 \end{bmatrix} \end{array}$$





First system of equations in (2) can be written as

$$\partial_t \psi + \mathcal{L}_1 \psi = \mathsf{f}_1$$

where $\mathcal{L}_1 \psi = \sum_{k=1}^3 \tilde{\mathbf{A}}_k \partial_k \psi + \mathbf{C}_1 \psi$ with $\tilde{\mathbf{A}}_k$ as in previous example and $\mathbf{C}_1 = \begin{bmatrix} iM\mathbf{I} & 0\\ 0 & -iM\mathbf{I} \end{bmatrix}$, $\mathbf{f}_1 = i\gamma^0 \psi \phi$.

For the second system of equations in (2) we introduce

$$\mathbf{v} = \begin{bmatrix} v_1 \\ v_2 \\ \mathbf{v}_3 \end{bmatrix} = \begin{bmatrix} \phi \\ \partial_t \phi \\ -\nabla \phi \end{bmatrix}$$

in order to get an evolution Friedrichs system

$$\partial_t v + \mathcal{L}_2 v = f_2,$$

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Using block diagonal structure, we get evolution Friedrichs system

$$\partial_t \mathbf{u} + \mathcal{L} \mathbf{u} = \mathbf{f},$$

where
$$\mathbf{u} = \begin{bmatrix} \psi & \mathbf{v} \end{bmatrix}^{\top}$$
, \mathcal{L} is Friedrichs operator with $\mathbf{A}_{k} = \begin{bmatrix} \tilde{\mathbf{A}}_{k} & 0\\ 0 & \bar{\mathbf{A}}_{k} \end{bmatrix}$,
 $\mathbf{C} = \begin{bmatrix} \mathbf{C}_{1} & 0\\ 0 & \mathbf{C}_{2} \end{bmatrix}$ and $\mathbf{f} = \begin{bmatrix} \mathbf{f}_{1} & \mathbf{f}_{2} \end{bmatrix}^{\top}$. Spaces involved:
 $L = \mathbf{H}^{2}(\mathbf{R}^{3}; \mathbf{C}^{9}),$
 $W = \{\psi \in \mathbf{H}^{2}(\mathbf{R}^{3}; \mathbf{C}^{9}) : \mathcal{L}_{1}\psi \in \mathbf{H}^{2}(\mathbf{R}^{3}; \mathbf{C}^{4})\} \times \mathbf{H}^{2}(\mathbf{R}^{3}) \times \mathbf{H}^{2}_{\mathrm{div}}(\mathbf{R}^{3}; \mathbf{C}^{3}).$
Moreover, \mathbf{f} is Lipshitz on $\mathbf{H}^{2}(\mathbf{R}^{3}; \mathbf{C}^{9})$ and D is trivial.

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Dirac-Maxwell system

$$\begin{cases} -\frac{i}{2\pi}(\gamma^{0}\partial_{t} + \gamma^{1}\partial_{1} + \gamma^{2}\partial_{2} + \gamma^{3}\partial_{3})\psi + m\beta\psi = \sum_{k=0}^{3}\mathcal{A}_{k}\gamma^{k}\psi \\ (-\frac{\partial^{2}}{\partial t} + \Delta)\mathcal{A}_{k} = -\gamma^{k}\psi\cdot\psi, \quad k = 0..3, \end{cases}$$
(3)

where $\gamma^0 = \mathbf{I}$ and γ^k , k = 1, 2, 3 as before. Unknown functions are $\psi : \mathbf{R}^{1+3} \to \mathbf{C}^4$ and $\mathcal{A} = \begin{bmatrix} \mathcal{A}_0 & \mathcal{A}_1 & \mathcal{A}_2 & \mathcal{A}_3 \end{bmatrix}^\top$, while $m \ge 0$ and $\beta = \begin{bmatrix} \mathbf{I} & 0 \\ 0 & -\mathbf{I} \end{bmatrix}$. Analog procedure as in previous example gives us Friedrichs system

$$\partial_{t} \mathbf{u} + \mathcal{L} \mathbf{u} = \mathbf{F},$$
where $\mathbf{u} = \begin{bmatrix} \psi \\ \mathbf{v}_{0} \\ \mathbf{v}_{1} \\ \mathbf{v}_{2} \\ \mathbf{v}_{3} \end{bmatrix}$, $\mathbf{v}_{k} = \begin{bmatrix} \mathcal{A}_{k} \\ \partial_{t} \mathcal{A}_{k} \\ -\nabla \mathcal{A}_{k} \end{bmatrix}$, $\mathbf{f}_{k} = \begin{bmatrix} 0 \\ \gamma^{k} \psi \cdot \psi \\ 0 \\ 0 \\ 0 \end{bmatrix}$, $\mathbf{k} = 0..3$
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Operator
$$\mathcal{L}$$
 is Friedrichs operator with $\mathbf{C} = \begin{bmatrix} \tilde{\mathbf{C}} & 0 \\ 0 & 0 \end{bmatrix} \in M_{24}(\mathbf{C}),$
 $\tilde{\mathbf{C}} = \begin{bmatrix} 2\pi im & 0 \\ 0 & -2\pi im \end{bmatrix}$ and
 $\mathbf{A}_{k} = \begin{bmatrix} \tilde{\mathbf{A}}_{k} & 0 & 0 & 0 & 0 \\ 0 & \bar{\mathbf{A}}_{k} & 0 & 0 & 0 \\ 0 & 0 & 0 & \bar{\mathbf{A}}_{k} & 0 \\ 0 & 0 & 0 & 0 & \bar{\mathbf{A}}_{k} \end{bmatrix}, \quad k = 1, 2, 3$
 $\mathbf{F} = \begin{bmatrix} \sum_{k=0}^{3} \mathcal{A}_{k} \gamma^{k} \psi & \mathbf{f}_{0} & \mathbf{f}_{1} & \mathbf{f}_{2} & \mathbf{f}_{3} \end{bmatrix}^{\mathsf{T}}$. Spaces involved:
 $L = \mathrm{H}^{2}(\mathbf{R}^{3}; \mathbf{C}^{24}),$

 $W = \{\psi \in \mathrm{H}^2(\mathbf{R}^3) : \bar{A}_1 \partial_1 \psi + \bar{A}_2 \partial_2 \psi + \bar{A}_3 \partial_3 \psi \in \mathrm{H}^2(\mathbf{R}^3) \} \times [\mathrm{H}^2(\mathbf{R}^3) \times \mathrm{H}^3(\mathbf{R}^3) \times \mathrm{H}^2_{\mathrm{div}}(\mathbf{R}^3)]^4.$





Shank you for your attention!